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I. Petkov, Vesselin II. Lazarov, Raitcho  
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V Petkov and R Lazarov (Editors)

Institute of Mathematics, Bulgarian Academy of Sciences

# Integral Equations and Inverse Problems

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D.G. VASIL'EV

## One can hear the dimension of a connected fractal in $R^2$

Let  $\Omega$  be a bounded domain in  $R^2$  with boundary  $\Gamma$ ; it is assumed that  $\Gamma$  has zero measure. We consider an eigenvalue problem

$$-\Delta u = \lambda u, \quad (1)$$

$$u|_{\Gamma} = 0, \quad (2)$$

where  $\Delta = \partial^2/\partial x_1^2 + \partial^2/\partial x_2^2$  is the Laplacian acting on  $\Omega$ . Problem (1), (2) is posed in variational form on functions from  $H_0^1(\Omega)$ .

Problem (1), (2) has a discrete spectrum  $\lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_k \leq \dots$ . By  $N(\lambda)$  we shall denote the eigenvalue distribution function, i.e. the number of eigenvalues  $\lambda_k$  below a given  $\lambda$ .

It is known [1] that the following asymptotics holds

$$N(\lambda) = \frac{S}{4\pi} \lambda + o(\lambda), \quad \lambda \rightarrow +\infty, \quad (3)$$

where  $S = |\Omega|$  is the area (Lebesgue measure) of  $\Omega$ . Thus the area of  $\Omega$  is uniquely recovered from the spectrum. The aim of this work is to show that a certain real number  $1 \leq d^{(i)} \leq 2$  having the sense of dimension and characterizing the degree of 'hairiness' of the boundary from the interior is also uniquely reconstructed from the spectrum. Such 'hairy' objects are called fractals [2].

Let us give some definitions.

By  $\Gamma_\varepsilon$  we shall denote the  $\varepsilon$ -neighbourhood of set  $\Gamma$ . The number  $d$  defined as the infimum of all positive  $\delta$  such that

$$\limsup_{\varepsilon \rightarrow +0} \varepsilon^{\delta-2} |\Gamma_\varepsilon| < +\infty \quad (4)$$

is called [3], [4] the Bouligand-Minkowski dimension of  $\Gamma$ .

By  $\Gamma_\varepsilon^{(i)}$ ,  $\Gamma_\varepsilon^{(e)}$  we shall denote subsets of  $\Gamma_\varepsilon$  determined by relations  $x \in \Omega$ ,  $x \notin \Omega$  respectively. Substituting in formula (4)  $\Gamma_\varepsilon$  by  $\Gamma_\varepsilon^{(i)}$ ,  $\Gamma_\varepsilon^{(e)}$  we define by analogy with  $d$  numbers  $d^{(i)}$ ,  $d^{(e)}$ . We will call these numbers respectively the interior and exterior Bouligand-Minkowski dimension of the boundary.

Finally by  $h$  we shall denote the Hausdorff dimension of the boundary defined [2] as the infimum of positive  $\eta$  such that

$$\lim_{\varepsilon \rightarrow +0} \left( \inf_{j \in J} \sum_j r_j^\eta \right) = 0. \quad (5)$$

The infimum in (5) is taken over all coverings of set  $\Gamma$  by balls of radii  $r_j < \varepsilon$ .

It can be shown (see also [3], [4]) that

$$d, d^{(i)}, d^{(e)}, h \in [1, 2], \quad d = \max(d^{(i)}, d^{(e)}), \quad h \leq d. \quad (6)$$

Let us introduce the function

$$Z(t) = \sum_{k=1}^{+\infty} \exp(-\lambda_k t)$$

$t > 0$ . It is known [5] that  $Z(t) < S/4\pi t$ .

We will call a compact set in  $R^2$  connected if it cannot be divided into two non-empty subsets with a smooth compact curve separating them.

The principal result of this work is contained in the following

**Theorem.** If  $\Gamma$  has only a finite number of connected components then

$$-2 \liminf_{t \rightarrow +0} \frac{\ln(S/4\pi t - Z(t))}{\ln t} = d^{(i)}. \quad (7)$$

Let us give an outline of the theorem's proof. Let  $u(x, y, t)$  denote the solution of the heat equation  $+\Delta u = \partial u/\partial t$  with boundary condition (2) and initial condition  $u|_{t=0} = \delta(x-y)$ ,  $y \in \Omega$ . Let  $u_0(x, y, t) = (4\pi t)^{-1} \exp(-|x-y|^2/4t)$  denote the fundamental solution of the heat equation. Of course

$$Z(t) = \int_{\Omega} u(x, x, t) dx, \quad (8)$$

$$0 \leq u(x, y, t) < u_0(x, y, t), \quad x, y \in \Omega, \quad t > 0. \quad (9)$$

Let  $\rho(x, \Gamma)$  denote the distance from  $x$  to  $\Gamma$ . It can be shown that for sufficiently small  $t$

$$u_0(x, x, t) - u(x, x, t) \leq c_{\kappa\mu} t^\mu, \quad \rho(x, \Gamma) \geq t^\kappa \quad (\forall \kappa < 1/2, \forall \mu > 0), \quad (10)$$

$$u_0(x, x, t) - u(x, x, t) > C t^{-1}, \quad \rho(x, \Gamma) \leq t^{1/2}. \quad (11)$$

In (11)  $C$  is a universal positive constant independent of  $\Omega$  and  $\Gamma' \neq \emptyset$  is a connected component of  $\Gamma$  with dimension  $d^{(i)}$ . Substituting (9)-(11) into (8) and using the definition of dimension  $d^{(i)}$  we obtain (7).

The most delicate point in the arguments outlined above is obtaining estimate (11). It is interesting that in the case  $\Omega \subset R^n$ ,  $n > 3$ , an analogous estimate  $u_0 - u \geq C t^{-n/2}$  does not hold and in place of (7) we have

$$-2 \liminf_{t \rightarrow +0} \frac{\ln(|\Omega|/(4\pi t)^{n/2} - Z(t))}{\ln t} \leq d^{(i)} \quad (12)$$

with effective examples of strict inequality. We must note that a somewhat weaker version of inequality (12) with  $d$  instead of  $d^{(i)}$  (see also (6)) in the right-hand side was obtained in [4].

In conclusion let us discuss Berry's conjecture [6] on the existence of a two-term asymptotic expansion

$$N(\lambda) = a\lambda^{n/2} + b\lambda^{v/2} + o(\lambda^{v/2}), \quad \lambda \rightarrow +\infty, \quad (13)$$

$a = (2\pi)^{-n} B_n |\Omega|$ ,  $B_n$  is the volume of the unit ball in  $R^n$ ,  $v$  is the 'dimension' of  $\Gamma$ ,  $b$  is a certain constant characterizing the  $v$ -dimensional 'volume' of  $\Gamma$ .

The first problem arising in connection with (13) is how to define dimension  $v$ ? Berry proposed  $v = h$  but a counterexample was constructed in [3]. It was proposed [3], [4] to take  $v = d$ . But it follows from (7), (12) that for the general case  $v = d$  also is not good because one can easily construct examples of domains  $\Omega$  with  $d^{(i)} < d$  (in this situation a contradiction can be avoided only if  $b = 0$  which is rather an unnatural assumption).

Another remark on Berry's conjecture is that the form  $\text{const} \cdot \lambda^{v/2}$  of the second asymptotic term is not obvious. If one considers a domain  $\Omega$  with a self-similar [2] boundary  $\Gamma$  one can expect the two-term  $N(\lambda)$  asymptotics to have the form

$$N(\lambda) = a\lambda^{n/2} + f(\ln \lambda)\lambda^{v/2} + o(\lambda^{v/2}),$$

where  $f(\cdot)$  is a periodic function with period  $2 \ln s$ ,  $s > 1$  is the linear similarity coefficient,  $v = d = d^{(i)} = d^{(e)}$ .

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D.G. Vasil'ev  
 Institute for Problems in Mechanics,  
 Academy of Sciences of USSR,  
 Prospect Vernadskogo 101,  
 117526 Moscow  
 USSR